

Gravitational waves from first-order phase transitions

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Work in progress.

With *C. Grojean*

In this work:

We determine the type of phase transitions which could be detected by LISA, LIGO, BBO in terms of the amount of supercooling and the duration of the phase transition.

BBO: Big Bang Observer,
second generation of space-interferometers

We point out that the characteristic frequency of the GW spectrum due to phase transitions in the temperature range $100 \text{ GeV} - 10^7 \text{ GeV}$ is in the window that will be probed by BBO.

initial motivation:

Is the electroweak phase transition 1st or 2nd order?

LHC might answer as it will shed light on the Higgs sector

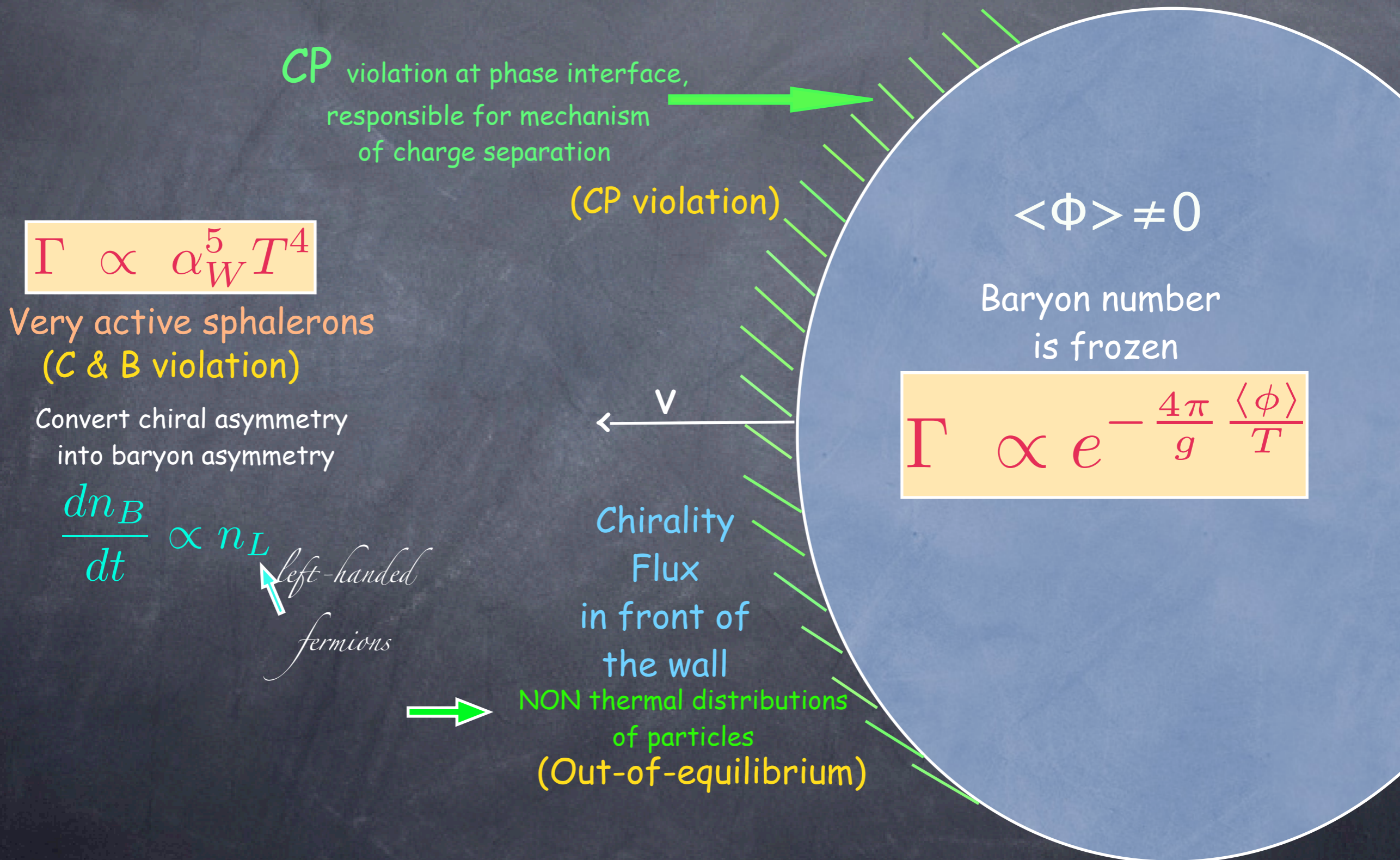
Question intensively studied within the Minimal Supersymmetric Standard Model (MSSM).

However, not so beyond the MSSM:

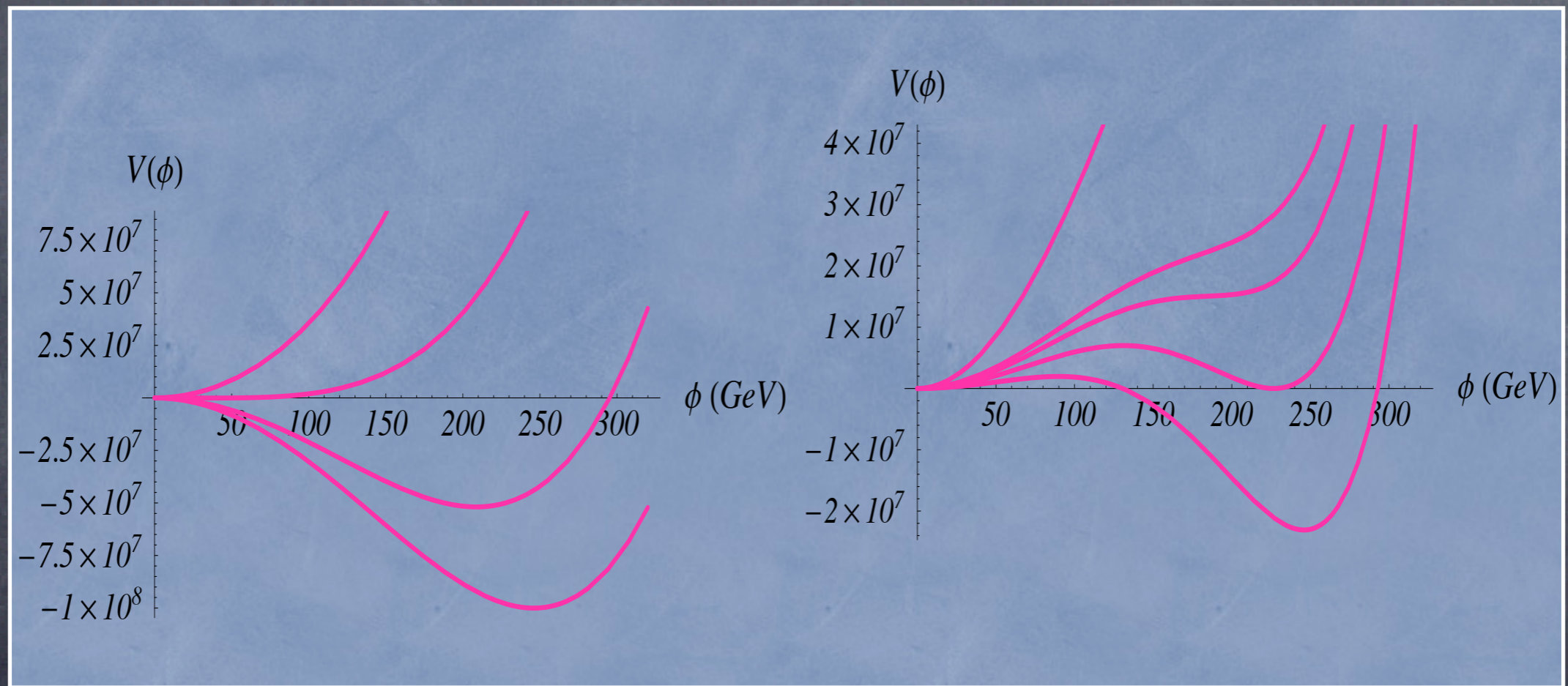
It is timely to investigate whether the recently proposed new models of Electroweak symmetry breaking (gauge-higgs unification in extra dimensions, composite Higgs, Little Higgs...) can lead to a first order electroweak phase transition.

Electroweak baryogenesis :

A beautiful mechanism to generate the matter-antimatter asymmetry of the universe involving EW physics only.



2nd order versus 1st order



Gravitational wave background from a strong first order PT

Kosowsky, Turner, Watkins'92
Kamionkowski, Kosowsky Turner '94

2 contributions: -from bubble collisions
-from turbulent fluid motions

Kosowsky, Mack, Kahniashvili'02
Dolgov, Grasso, Nicolis'02

2 basic quantities determine the GW background generated during a 1st order PT

⇒ α : ratio of energy density of false vacuum to thermal energy density of the plasma

⇒ β : rate of variation of nucleation rate Γ ; ($\Gamma = \Gamma_0 e^{-\beta t}$)

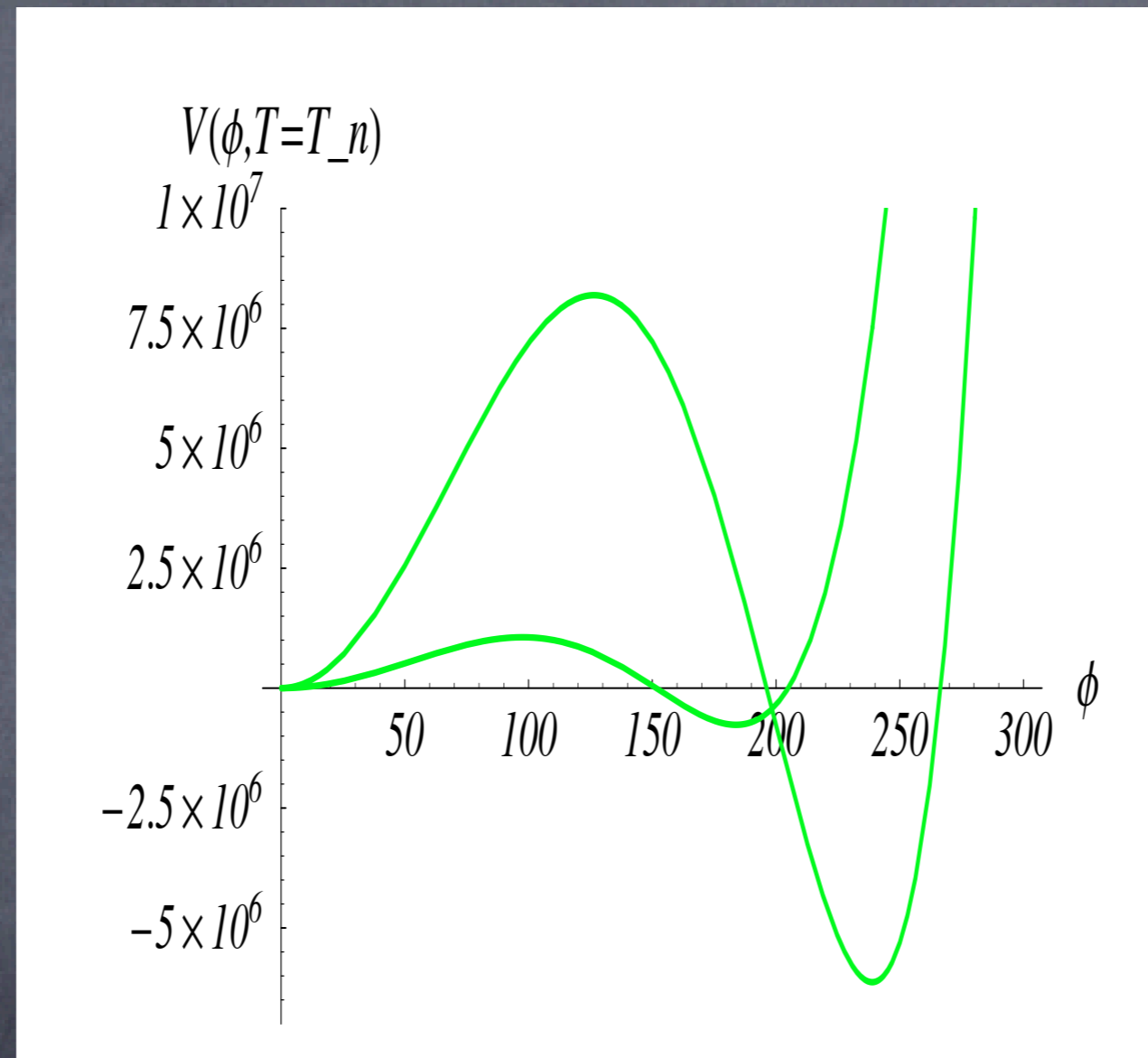
$\beta^{-1} = \Gamma / \dot{\Gamma}$: ~duration of PT

$2\pi f \approx \beta$ characteristic frequency of GW @ time of production

The stronger is the transition, the larger is α and the smaller is β (the larger amount of supercooling)

α and β : entirely determined by the effective scalar potential at high temperature

Shape of the potential at the nucleation temperature



× Frequency of the signal

$$f_{\text{coll}} \approx 5.2 \times 10^{-3} \text{mHz} \frac{\beta}{H_*} \frac{T_*}{100 \text{GeV}} \left[\frac{g_*}{100} \right]^{1/6}$$
$$f_{\text{turb}} \approx 0.65 \frac{u_s}{v_b} f_{\text{coll}}$$

from, e.g, A. Nicolis '03

× Amplitude of the signal at the peak frequency

$$\Omega_{\text{gw}} h^2 \propto \left[\frac{H_*}{\beta} \right]^2 \times f(\alpha)$$

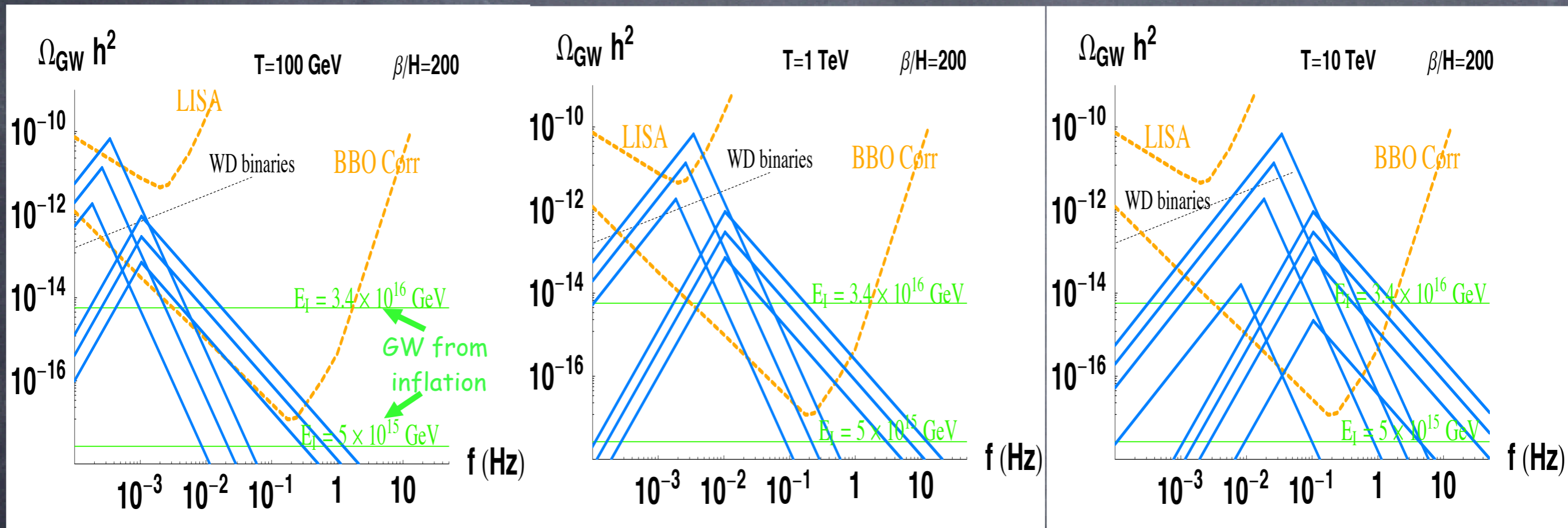
where $f(\alpha)$ increases with α

The **gravitational wave spectrum** from a 1st order **electroweak** PT is peaked around the **milliHertz** frequency, just in the band of the planned space interferometer LISA.

Spectrum of gravitational waves produced at 1st order phase transitions

1st blue peak: GW from bubble collisions

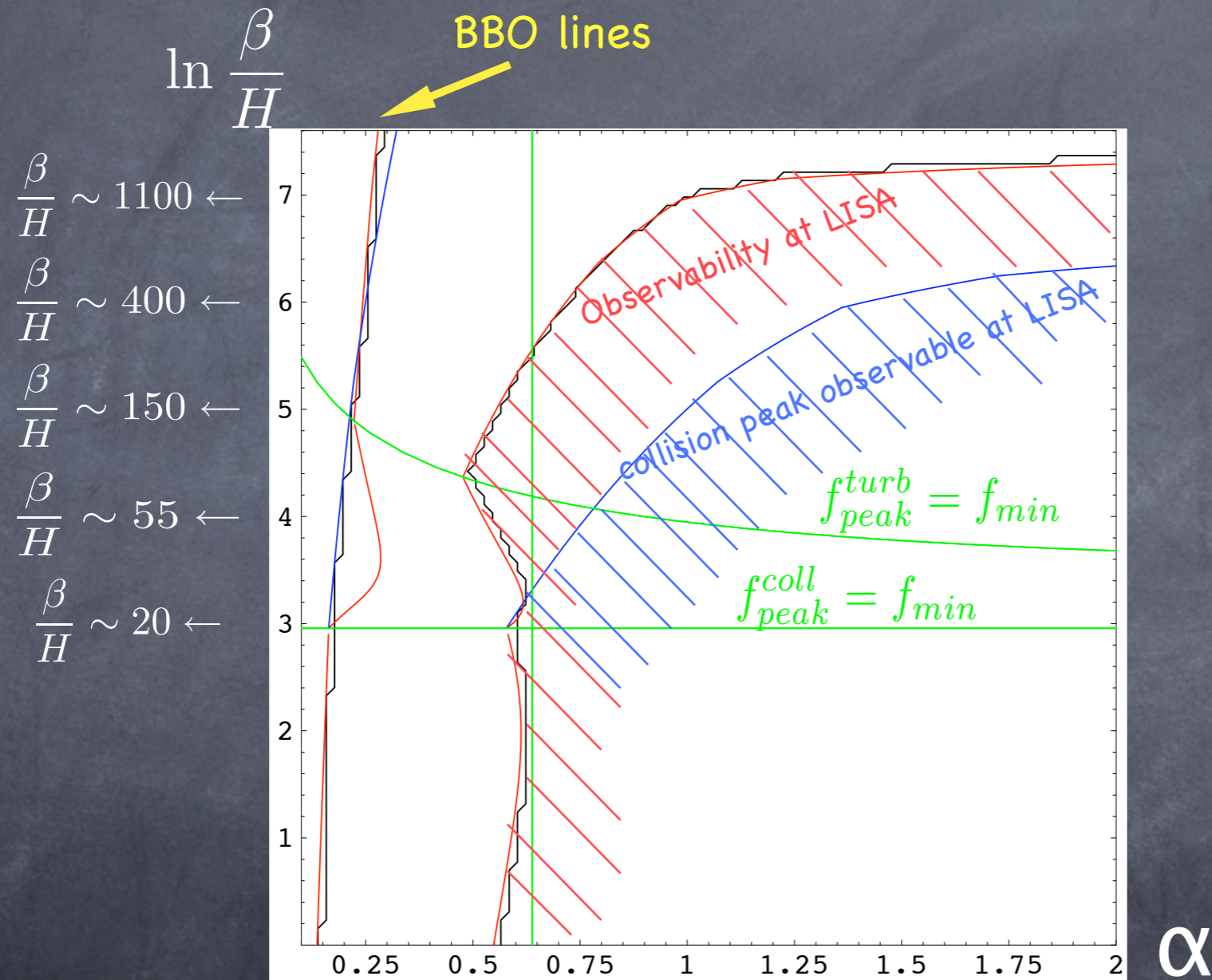
2nd blue peak: GW from turbulence



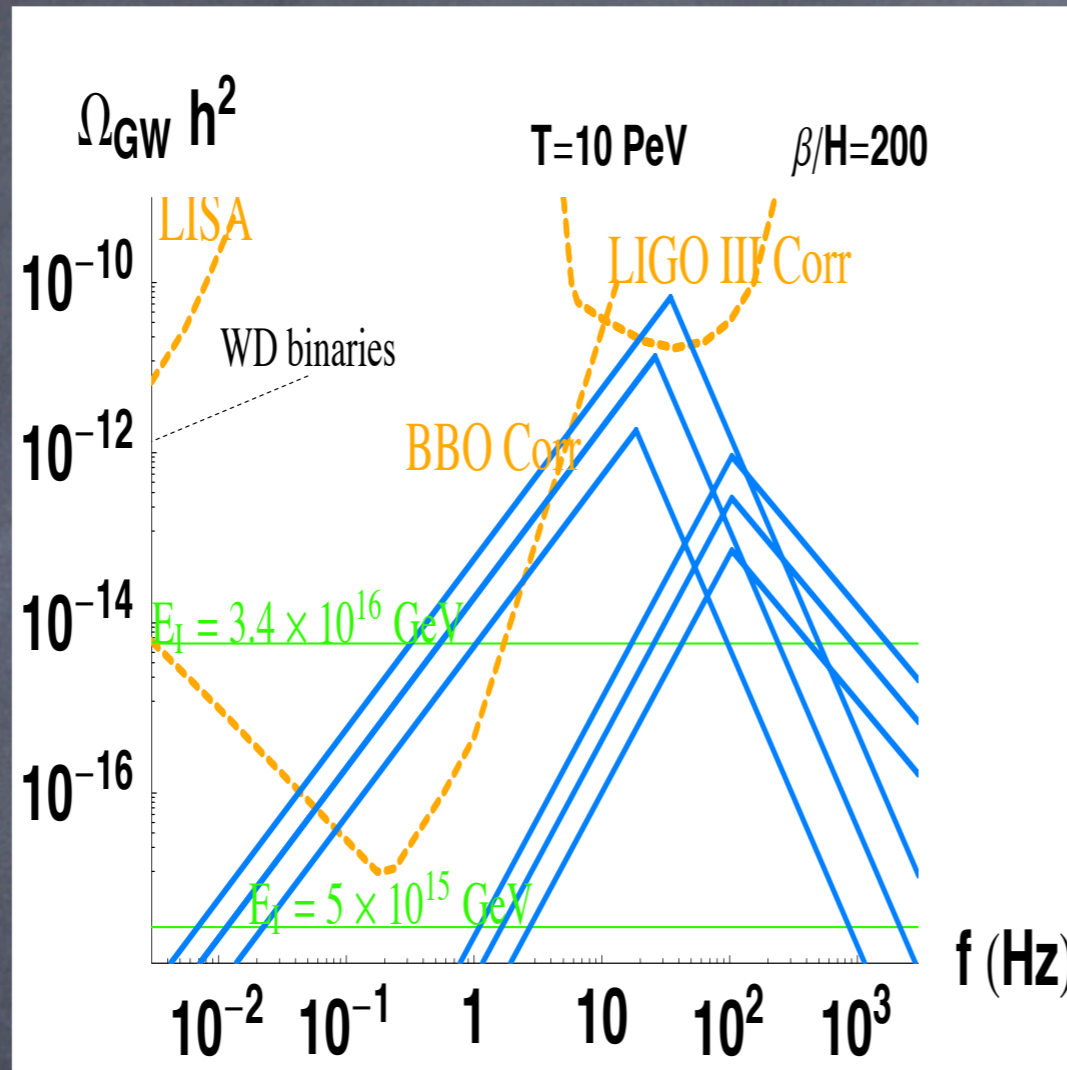
$$\alpha = 0.3, 0.5, 0.8$$

Observability of the EWPT at LISA and BBO in the α - β/H plane

$T=100$ GeV



A phase transition at $T=10^7$ GeV could be observed both at LIGO and BBO:



$\alpha = 0.3, 0.5, 0.8$

Summary

- LISA, LIGO and BBO: can probe phase transitions in the temperature range $100 \text{ GeV} - 10^7 \text{ GeV}$
- The GW spectrum is a function of 2 quantities: α (latent heat) and β^{-1} (duration of PT) which can be computed for any given scalar potential.
- For each T , we determined the region in the (α, β) plane leading to an observable signal. Observation or non-observation of GW will put constraints on the parameters of these potentials.
- At LHC, the higgs mass is measured but not the quartic or cubic Higgs self coupling. A linear collider can provide this information, which timescale is even beyond LISA. LISA could start constraining parameters of the Higgs sector before a linear collider...
- The signal from PT at $T \sim 1-100 \text{ TeV}$ could entirely screen the relic GW background expected from inflation.